

Theoretical calculation and experimental measurement on production cross section of Z boson in μ-μ channel with Monte Carlo simulated data

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Physics (Intensive)

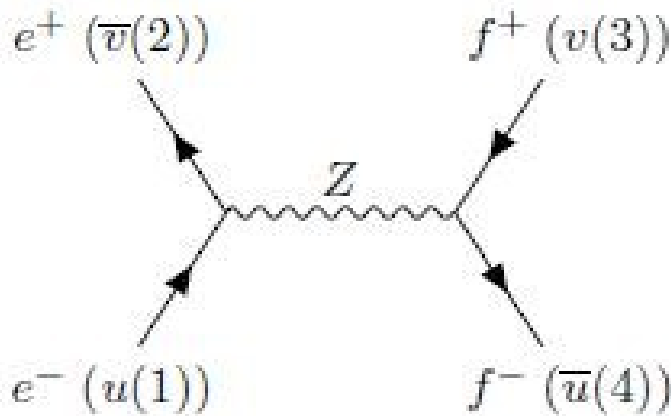
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Abstract

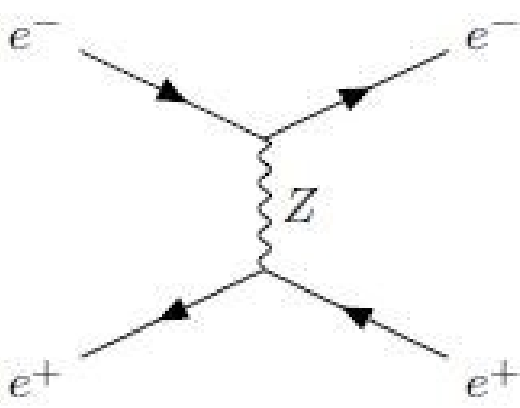
Basics of Quantum field theory was discussed in this research. We proposed a way of calculating the production cross section of Z-boson in tree-level Feynman Diagrams using the existing Dirac Feynman Rules. Results are then compared with colliding data in Monte Carlo Event generator available in PYTHIA 8. Consistent results are attained.

Methods

The aim for this section is to find out the cross-section for the process $e^+e^- \rightarrow Z \rightarrow e^+e^-$ and $\rightarrow \mu^+\mu^-$. We start from the Feynman diagrams.



There is another possible Feynman diagram for $e^+e^- \rightarrow Z \rightarrow e^+e^-$ but not $\mu^+\mu^-$.



Obviously, these processes are mediated by the weak force carrier Z boson, thus we have to use the corresponding weak neutral current vertex factor and propagator.

To simplify things up, we will neglect the mass of electrons and muons, because otherwise this process will have a much smaller cross-section compared to QED processes. This simplification is justified as Z boson is much heavier than electrons and muons.

On the other hand, we will consider the Z vertex factor:

$$\frac{-ig_Z}{2}\gamma^\mu(c_V^f - c_A^f\gamma^5)$$

from now on we take $c_V - c_A\gamma^5 = \Gamma$, so new form will be

$$\frac{-ig_Z}{2}\gamma^\mu\Gamma$$

We first find out the matrix element for the s-channel diagram. By Feynman rules we have:

$$\int \frac{d^4q}{(2\pi)^4}(2\pi)^8\delta(q-p_1-p_2)\delta(q-p_3-p_4)\overline{u}(4)(\frac{-ig}{2}\gamma^\mu\Gamma)v(3)\left[\frac{-i(g_{\mu\nu}-q_\mu q_\nu/M^2)}{q^2-M^2}\right]\overline{v}(2)(\frac{-ig}{2}\gamma^\mu\Gamma)u(1)$$
$$\mathcal{M}_s=\frac{g_Z^2}{4}[\overline{u}(4)\gamma^\mu\Gamma v(3)]\frac{g_{\mu\nu}-q_\mu q_\nu/M^2}{q^2-M^2}[\overline{v}(2)\gamma^\nu\Gamma u(1)]$$

Note that the term with the factor $q_\mu q_\nu/M^2$ vanishes if we neglect the mass of electron and muon, as for example,

$$\overline{u}(4)\not{q}\Gamma v(3)=\overline{u}(4)\not{p}_4X+X'\not{p}_3v(3)=0$$

by the massless Dirac equation $\gamma^\mu p_\mu v=0$ and its adjoint $\overline{u}\gamma^\mu p_\mu=0$. (X,X' are some finite terms.)

Hence, the norm square of \mathcal{M}_s is:

$$|\mathcal{M}_s|^2=\frac{g_Z^4}{16(q^2-M^2)^2}[\overline{u}(4)\gamma^\mu\Gamma v(3)][\overline{v}(2)\gamma_\nu\Gamma u(1)][\overline{u}(4)\gamma^\mu\Gamma v(3)]^*[\overline{v}(2)\gamma_\nu\Gamma u(1)]^*$$

where the asterisk represent complex conjugate.

Typically, if the process has specific spins and polarization configurations, we can insert the appropriate spinors and/or ϵ to work out \mathcal{M} and $|\mathcal{M}|^2$. But this is more like a subset of a real world experiment: specific spins are not favoured over any other configuration—we just do experiment with a beam of initial particles that have random spins. Therefore, what we can do is to compute the expected value of $|\mathcal{M}|^2$ which is by averaging over all initial spins and summing all final spins .

To proceed, we will employ the Casimir’s Trick to reduce the complexity of the calculations:

$$\sum_{\text{all spin}}[\overline{u}(a)\Gamma_1u(b)][\overline{u}(a)\Gamma_2u(b)]^*=Tr[\Gamma_1(\not{p}_b+m_b c)\Gamma_2(\not{p}_a+m_a c)]$$

It is a very useful tool since we can get rid of the spinors, and once we do the summation over spins, the computations are trim down into trace identities of some γ matrices.

So, for electron-positron pair scattering via Z boson at high energy, we have:

$$\langle|\mathcal{M}|^2\rangle=\langle|\mathcal{M}_1-\mathcal{M}_2|^2\rangle=\langle|\mathcal{M}_1|^2\rangle-2Re(\langle\mathcal{M}_1^*\mathcal{M}_2\rangle)+\langle|\mathcal{M}_2|^2\rangle$$
$$\langle|\mathcal{M}|^2\rangle=\frac{1}{2}\frac{g^4}{(q_1^2-M^2)^2}\{(c_V^2+c_A^2)^2[(p_1\cdot p_3)(p_2\cdot p_4)+(p_1\cdot p_4)(p_2\cdot p_3)]+4c_V^2c_A^2[(p_1\cdot p_3)(p_2\cdot p_4)-(p_1\cdot p_4)(p_2\cdot p_3)]\}$$
$$+\frac{1}{2}\frac{g^4}{(q_2^2-M^2)^2}\{(c_V^2+c_A^2)^2[(p_1\cdot p_3)(p_2\cdot p_4)+(p_3\cdot p_4)(p_2\cdot p_1)]+4c_V^2c_A^2[(p_1\cdot p_3)(p_2\cdot p_4)-(p_3\cdot p_4)(p_2\cdot p_1)]\}$$
$$+\frac{g^4}{2(q_1^2-M^2)(q_2^2-M_2)}(c_V^4+6c_V^2c_A^2+c_A^4)(p_2\cdot p_4)(p_1\cdot p_3)$$
$$\langle|\mathcal{M}|^2\rangle=(\frac{g^2E^2}{(2E)^2-M^2})^2[(c_V^2+c_A^2)^2(1+\cos^2\theta)-8c_V^2c_A^2\cos\theta]$$
$$+\frac{g^4}{2(4E^2\sin^2(\theta/2)+M^2)^2}[(c_V^2+c_A^2)^2(E^4(1-\cos\theta)^2+4E^4)+4c_V^2c_A^2(E^4(1-\cos\theta)^2-4E^4)]$$
$$-\frac{g^4(c_V^4+6c_V^2c_A^2+c_A^4)}{2((2E)^2-M^2)(4E^2\sin^2(\theta/2)+M^2)}(E^4(1-\cos\theta)^2)$$

Theoretical results

We give the cross section for $e^+e^- \rightarrow \mu^+\mu^-$. The calculation for $e^+e^- \rightarrow \mu^+\mu^-$ is similar. Suppose CM frame energy $2E=500GeV$. Neglect electron and muon mass. Neglect the running constant problem and simply take $g_Z=0.718$, $c_A=-1/2$, $c_V=-1/2+2\sin^2\theta_W=-0.0372$ so $c_A^2+c_V^2\approx 1/4$. Take $M=91.2GeV$. From Golden Rule we can derive the cross section:

$$\frac{d\sigma}{d\Omega}=(\frac{1}{8\pi})^2\frac{|p_f|}{|p_i|}\frac{|\mathcal{M}|^2}{(2E)^2}$$
$$\sigma=\frac{1}{(16\pi E)^2}\int d\Omega|\mathcal{M}|^2$$
$$|\mathcal{M}|^2=(\frac{g^2E^2}{(2E)^2-M^2})^2((c_V^2+c_A^2)^2(1+\cos^2\theta)-8c_V^2c_A^2\cos\theta)$$
$$\approx(\frac{g^2E^2}{(2E)^2-M^2})^2((1+\cos^2\theta)/16-8c_V^2c_A^2\cos\theta)$$
$$\sigma=(\frac{g^2E}{16\pi((2E)^2-M^2)})^22\pi\int_0^\pi\frac{1+\cos^2\theta}{16}-8c_V^2c_A^2\cos\theta$$
$$=(\frac{g^2E}{16\pi((2E)^2-M^2)})^2\frac{\pi}{8}(\frac{3}{2}\pi)$$
$$=\frac{3g^4E^2}{16^3(4E^2-M^2)^2}$$
$$\approx 8.33\times 10^{-10}GeV^{-2}$$

which equals a cross section of 0.324 pico-barn.

To turn to MCEG and experimental data, we use PYTHIA8:

```
#include "Pythia8/Pythia.h"
using namespace Pythia8;

int main(int argc, char* argv[]) {
    Pythia pythia;
    pythia.readFile(argv[1]);

    pythia.init();

    int n;
    for(int i=0; i<50000;++i){
        pythia.next();
        for(int j=0; j<pythia.event.size();++j){
            if(pythia.event[j].id() == 11 || pythia.event[j].id() == -11){
                n++;
            }
        }
        pythia.stat();
        cout << "##### " <<endl;
        return 0;
    }
```

Beams:idA = 11
Beams:idB = -11
Beams:eCM = 500
WeakZ0:gmZmode = 2
WeakBosonExchange:ff2ff(t:gmZ) = on
Next:numberShowEvent = 8
23:onIfAny = 13

After ~50000 events, we have filtered the useful information and store them in an out file:

no	id	name	status	mothers	daughters	colours	p_X	p_Y	p_Z	e	m
0	90	(system)	-11	0	0	0	0.000	0.000	0.000	500.000	500.000
1	11	(e-)	-12	0	0	21	0	0	0.000	250.000	0.000
2	-11	(e+)	-12	0	0	22	0	0	0.000	-250.000	0.000
3	11	(e-)	-21	6	6	5	0	0	0.000	249.902	249.902
4	-11	(e+)	-21	7	0	5	0	0	0.000	-8.790	8.790
5	23	(Z0)	-22	3	4	8	8	0	0.000	241.112	93.739
6	11	(e-)	-42	10	0	3	3	0	-0.000	249.902	0.000
7	-11	(e+)	-41	11	11	9	4	0	0.000	-249.804	249.804
8	23	(Z0)	-44	5	5	12	12	0	-16.867	238.898	258.692
9	22	(gamma)	-43	7	0	13	13	0	16.867	-238.799	241.014
10	11	(e-)	-41	15	15	14	6	0	0.000	0.000	250.000
11	-11	(e+)	-42	16	0	7	7	0	-0.000	0.000	-249.804
12	23	(Z0)	-44	8	8	17	17	0	-16.868	27.895	238.898
13	22	(gamma)	-44	9	9	18	18	0	16.867	-27.892	-238.799
14	22	(gamma)	-43	10	0	19	19	0	0.001	-0.003	0.098
15	11	(e-)	-42	21	21	10	10	0	-0.000	0.000	250.000
16	-11	(e+)	-41	22	22	20	11	0	0.000	-0.000	-250.000
17	23	(Z0)	-44	12	12	23	23	0	-16.868	27.895	238.897
18	22	(gamma)	-44	13	13	24	24	0	16.868	-27.893	-238.799
19	22	(gamma)	-44	14	14	25	25	0	0.001	-0.003	0.098
20	22	(gamma)	-43	16	0	26	26	0	-0.002	0.001	-0.196
21	11	(e-)	-61	1	0	15	15	0	-0.000	0.000	250.000
22	-11	(e+)	-61	2	0	16	16	0	0.000	0.000	-250.000
23	23	(Z0)	-62	17	17	29	30	0	-16.868	27.895	238.897
24	22	gamma	-62	18	18	0	0	0	16.868	-27.893	-238.799
25	22	gamma	-62	19	19	0	0	0	0.001	-0.003	0.098
26	22	gamma	-62	20	20	0	0	0	-0.002	0.001	-0.196
27	22	gamma	-63	1	0	0	0	0	0.000	0.000	0.000
28	22	gamma	-63	2	0	0	0	0	0.000	0.000	0.000
29	13	mu-	23	23	0	0	0	0	-25.613	46.095	231.110
30	-13	mu+	23	23	0	0	0	0	8.745	-18.200	7.787
				Charge sum: 0.000		Momentum sum:		-0.000	-0.000	500.000	500.000
----- End PYTHIA Event Listing -----											
----- PYTHIA Event and Cross Section Statistics -----											
Subprocess		Code	Number of events				sigma +- delta				
			Tried	Selected	Accepted	(estimated) (mb)					

f fbar -> gamma*/Z0		221	4474	500	500	4.030e-10		9.692e-12			
sum			4474	500	500	4.030e-10		9.692e-12			
----- End PYTHIA Event and Cross Section Statistics -----											

From the simulation we found out that the experimental cross section of the t-channel scattering of lepton pairs to Z boson is roughly 0.4 pb, in the same magnitude as the theoretical calculations.

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Together we have written an introductory QFT reports on the basics in calculating cross section of elementary tree-level processes.

An Introduction to Quantum Field Theory
Cross section calculation of Z⁰ boson production

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Bibilography

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